

# RWQF6

$$1) \Sigma^{23} = \frac{i}{4} [\gamma^2, \gamma^3] = \frac{i}{2} \gamma^2 \gamma^3 = \frac{1}{2} \mathbb{1} \otimes \sigma^1$$

where we used Dirac's representation for the  $\gamma$ 's

$$\gamma^j = i \sigma^2 \otimes \sigma^j \quad j = 1, 2, 3$$

similarly we have  $\Sigma^{31} = \frac{1}{2} \mathbb{1} \otimes \sigma^2$ ;  $\Sigma^{12} = \frac{1}{2} \mathbb{1} \otimes \sigma^3$

Thus we have

$$\begin{aligned} [S^j, S^k] &= \frac{1}{4} \mathbb{1} \otimes [\sigma^j, \sigma^k] \\ &= \frac{1}{4} \mathbb{1} \otimes \sigma^l \ 2i \varepsilon^{jkl} \\ &= \frac{i}{2} \varepsilon^{jkl} \left( \frac{1}{2} \mathbb{1} \otimes \sigma^l \right) \\ &= \frac{i}{2} \varepsilon^{jkl} S^l \end{aligned}$$

2) We need to check  $\partial_\mu \bar{\psi} \gamma^\mu \psi = 0$

$$\partial_\mu j^\mu = \partial_\mu \bar{\Psi} \gamma^\mu \Psi + \bar{\Psi} \gamma^\mu \partial_\mu \Psi$$

clearly we need to use Dirac's equation

$$i \gamma^\mu \partial_\mu \Psi - m \Psi = 0; \quad i \partial_\mu \bar{\Psi} \gamma^\mu + m \bar{\Psi} = 0$$

Thus we have:

$$\partial_\mu j^\mu = i m \bar{\Psi} \Psi - i m \bar{\Psi} \Psi = 0$$

③ In the case at hand we have:

$$i \frac{\partial \phi}{\partial t} = \frac{1}{2m} (-i \partial_j + q A^j) (-i \partial_j + q A^j) \phi + \left( \frac{q}{2m} \sigma^j B^j - q A^0 \right) \phi$$

where  $A^j = \frac{1}{2} \epsilon^{jnl} B^k x^l$ . Then the 1<sup>st</sup> line is

$$\begin{aligned} \text{1<sup>st</sup> line} = & \frac{1}{2m} \left[ -\partial_j \partial_j \phi - i q \phi (\partial_j A^j) + q^2 A^j A^j \phi \right. \\ & \left. - 2i q \frac{1}{2} \epsilon^{jnl} B^k x^l \partial_j \phi \right] \end{aligned}$$

Now  $\partial_j A^j = \frac{1}{2} \epsilon^{ijn} B^n = 0$

As  $\epsilon^{ijn}$  is zero when two indices are equal. We'll also neglect the  $A^j A^j$  term.

Then the Pauli equation reduces to

$$i \frac{\partial \phi}{\partial t} = -\frac{1}{2m} \partial_j \partial_j \phi + \frac{q}{2m} (\sigma^k - i \epsilon^{kln} x^l \partial_j) B^n \phi - q A^0 \phi$$

Now  $S^k = \frac{1}{2} \sigma^k$  and  $L^k = -i \epsilon^{kln} x^l \partial_j$

Thus we have

$$i \frac{\partial \phi}{\partial t} = -\frac{1}{2m} \partial_j \partial_j \phi + \frac{q}{2m} (2S^k + L^k) B^k \phi - q A^0 \phi$$

(4)  $N = \int d^3z \hat{\phi}^\dagger(z) \hat{\phi}(z)$

$$\hat{H} = \underbrace{\int d^3x \hat{\phi}^\dagger(x) \left[ -\frac{\hbar^2}{2m} \nabla^2 + V_1(x) \right] \hat{\phi}(x)}_{H_0} +$$

$$\frac{1}{2} \int d^3x d^3y \underbrace{\hat{\phi}^\dagger(x) \hat{\phi}^\dagger(y) V_2(x,y) \hat{\phi}(x) \hat{\phi}(y)}_{H_1}$$

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By using the commutation relations

$$[\hat{\phi}(x), \hat{\phi}^\dagger(y)] = \delta^3(x-y)$$

we get

$$[\hat{N}, \hat{\phi}(x)] = \int d^3z \hat{\phi}(z) [\hat{\phi}^\dagger(z), \hat{\phi}(x)] = -\hat{\phi}(x)$$

$$[\hat{N}, \hat{\phi}^\dagger(x)] = \int d^3z [\hat{\phi}(z), \hat{\phi}^\dagger(x)] \hat{\phi}^\dagger(z) = \hat{\phi}^\dagger(x)$$

Notice the sign difference! Then

$$[\hat{N}, \hat{\phi}(x) \hat{\phi}^\dagger(x)] = n \hat{\phi}(x) \hat{\phi}^\dagger(x) - n \hat{\phi}^\dagger(x) \hat{\phi}(x) = 0$$

for any positive integer  $n$ . Finally

$$[\hat{N}, \hat{\phi}^\dagger(x) \nabla^2 \hat{\phi}(x)] = \hat{\phi}^\dagger(x) \nabla^2 \hat{\phi}(x) + \hat{\phi}(x) [\hat{N}, \nabla^2 \hat{\phi}(x)]$$

The commutator in the last term is

$$-\int d^3z \hat{\phi}(z) \frac{\partial}{\partial x^i} \frac{\partial}{\partial x^j} \delta^3(z-x) = -\frac{\partial}{\partial x^i} \frac{\partial}{\partial x^j} \int d^3z \hat{\phi}(z) \delta^3(z-x)$$

$$= -\nabla^2 \hat{\phi}(x). \text{ Thus } [\hat{N}, \hat{H}_0 + \hat{H}_1] = 0$$